# Short Course in Quantum Information Lecture 2



# Formal Structure of Quantum Mechanics





# **Course Info**

All materials downloadable @ website

http://info.phys.unm.edu/~deutschgroup/DeutschClasses.html

# <u>Syllabus</u>

Lecture 1: Intro Lecture 2: Formal Structure of Quantum Mechanics Lecture 3: Entanglement Lecture 4: Qubits and Quantum Circuits Lecture 5: Algorithms Lecture 6: Error Correction Lecture 7: Physical Implementations Lecture 8: Quantum Cryptography





# **Lecture 1: Review**

- Information is *physical*: The ability of a machine to perform, e.g., computation is constrained by the laws of physics.
- Information: what we know.

•Bayesian probability assignments based on prior knowledge.

- Quantum theory has its own logical rules.
  - Assign complex amplitude to quantum process.
  - Add amplitudes for *indistinguishable* processes.
  - Absolute square of amplitude gives probability of finding outcome --> Interference of outcomes.
  - Measurement "collapses" state assignment.
  - No local realistic description of "hidden variables".





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- <u>Syllabus</u>
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# **The Ingredients of Quantum Theory**



#### Mathematical Structure (Hilbert space). Vector space over complex numbers with inner product.





# **Review: Real vector space (Euclidean)**

**Linear superposition**  $\mathbf{W} = c_1 \mathbf{V} + c_2 \mathbf{U}$ 





 $\underline{\mathsf{Inner product}} \qquad \|\mathbf{V}\|^2 \equiv \mathbf{V} \cdot \mathbf{V}$ 

$$\mathbf{Y} \cdot \mathbf{W} = \mathbf{Y} \cdot (c_1 \mathbf{V} + c_2 \mathbf{U}) = c_1 \mathbf{Y} \cdot \mathbf{V} + c_2 \mathbf{Y} \cdot \mathbf{U}$$

<u>Orthonormal basis</u>  $\mathbf{e}_x \cdot \mathbf{e}_x = \mathbf{e}_y \cdot \mathbf{e}_y = 1$   $\mathbf{e}_x \cdot \mathbf{e}_y = 0$ 

$$V_x = \mathbf{e}_x \cdot \mathbf{V} \qquad V_y = \mathbf{e}_y \cdot \mathbf{V} \qquad \|\mathbf{V}\|^2 = V_x^2 + V_y^2$$





# **Complex Vector Space: Analytic Signals**

Consider an oscillating electric vector,  $\mathbf{E}(t) = E_0 \cos(\omega t) \mathbf{e}_r$ 

**Intensity:** 
$$I = \frac{1}{T} \int_{0}^{T} |\mathbf{E}(t)|^{2} = \frac{1}{2} E_{0}^{2}$$

 $\mathbf{E}(t) = \operatorname{Re}(\tilde{\mathbf{E}}(t))$ **Complex signal:**  $\tilde{\mathbf{E}}(t) = E_0 \mathbf{e}_x e^{-i\omega t}$ 

$$\tilde{\mathbf{E}}^* \cdot \tilde{\mathbf{E}} = \left( E_0 \mathbf{e}_x \ e^{+i\omega t} \right) \cdot \left( E_0 \mathbf{e}_x \ e^{-i\omega t} \right) = E_0^2 \left( e^{+i\omega t} \ e^{-i\omega t} \right) = E_0^2 \blacksquare I = \frac{1}{2} \tilde{\mathbf{E}}^* \cdot \tilde{\mathbf{E}}$$

**Circular polarizaton:**  $\mathbf{E}(t) = E_0 (\cos(\omega t) \mathbf{e}_x + \sin(\omega t) \mathbf{e}_y) = \operatorname{Re}(\tilde{\mathbf{E}}(t))$ 

$$\tilde{\mathbf{E}}(t) = E_0 \Big( \mathbf{e}_x + i\mathbf{e}_y \Big) e^{-i\omega t} \qquad I = \frac{1}{2} \tilde{\mathbf{E}}^* \cdot \tilde{\mathbf{E}} = \frac{1}{2} E_0^2 (\mathbf{e}_x - i\mathbf{e}_y) \cdot (\mathbf{e}_x + i\mathbf{e}_y) = E_0^2$$
Normalized
polarization
$$\mathbf{e}_{\pm} \equiv \frac{\mathbf{e}_x \pm i\mathbf{e}_y}{\sqrt{2}} \qquad \left\| \mathbf{e}_{\pm} \right\|^2 = \mathbf{e}_{\pm}^* \cdot \mathbf{e}_{\pm} = 1 \qquad \mathbf{e}_{\pm}^* \cdot \mathbf{e}_{-} = 0$$

pola vector.

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### **Motivation: Polarization Optics**



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Malus's Law:  

$$\mathbf{E}_{V} = \mathbf{e}_{V} (\mathbf{e}_{V} \cdot \mathbf{E}_{0}) = \mathbf{e}_{V} (\mathbf{e}_{V} \cdot \vec{\varepsilon}) E_{0} = \mathbf{e}_{V} E_{0} \cos\theta$$

$$I_{V} = \frac{1}{2} \mathbf{E}^{*} \cdot \mathbf{E} = |\mathbf{e}_{V} \cdot \vec{\varepsilon}|^{2} \left(\frac{1}{2} |E_{0}|^{2}\right) = \cos^{2}\theta I_{0}$$



## Measurement in another basis (1)



Malus's Law:

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$$\mathbf{E}_{D\pm} = \mathbf{e}_{D\pm} \left( \mathbf{e}_{D\pm} \cdot \mathbf{E}_0 \right) = \mathbf{e}_{D\pm} \left( \mathbf{e}_{D\pm} \cdot \vec{\varepsilon} \right) E_0 = \mathbf{e}_{D\pm} \left( \frac{\cos\theta \pm \sin\theta}{\sqrt{2}} \right) E_0$$
$$I_{D\pm} = \frac{1}{2} \mathbf{E}_{D\pm}^* \cdot \mathbf{E}_{D\pm} = \left| \mathbf{e}_{D\pm} \cdot \vec{\varepsilon} \right|^2 \left( \frac{1}{2} \left| E_0 \right|^2 \right) = \left( \frac{\cos\theta \pm \sin\theta}{\sqrt{2}} \right)^2 I_0$$



# Measurement in another basis (2)



Malus's Law:

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$$\mathbf{E}_{R,L} = \mathbf{e}_{R,L} \left( \mathbf{e}_{R,L}^* \cdot \mathbf{E}_0 \right) = \mathbf{e}_{R,L} \left( \mathbf{e}_{R,L}^* \cdot \vec{\varepsilon} \right) E_0 = \mathbf{e}_{D\pm} \left( \frac{\cos\theta \mp i\sin\theta}{\sqrt{2}} \right) E_0 = \mathbf{e}_{D\pm} \left( \frac{e^{\mp i\theta}}{\sqrt{2}} \right) E_0$$
$$I_{R,L} = \frac{1}{2} \mathbf{E}_{R,L}^* \cdot \mathbf{E}_{R,L} = \left| \mathbf{e}_{R,L} \cdot \vec{\varepsilon} \right|^2 \left( \frac{1}{2} \left| E_0 \right|^2 \right) = \frac{1}{2} I_0$$



### **Repeated measurement**



Malus's Law:

 $\mathbf{E}'_{V} = \mathbf{e}_{V} \left( \mathbf{e}_{V} \cdot \mathbf{E}_{V} \right) = \mathbf{e}_{V} E_{V} \qquad \mathbf{E}'_{H} = \mathbf{e}_{H} \left( \mathbf{e}_{H} \cdot \mathbf{E}_{V} \right) = 0$  $I'_{V} = \frac{1}{2} \mathbf{E}'^{*}_{V} \cdot \mathbf{E}'_{V} = I_{V} \qquad I'_{H} = 0$ 

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### **Repeated measurement**



Malus's Law:

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$$\mathbf{E}'_{V} = \mathbf{e}_{V} \left( \mathbf{e}_{V} \cdot \mathbf{E}_{D+} \right) = \mathbf{e}_{V} \left( \mathbf{e}_{V} \cdot \mathbf{e}_{D+} \right) \left( \mathbf{e}_{D+} \cdot \mathbf{e}_{V} E_{V} \right) = \frac{1}{2} \mathbf{E}_{V} \Longrightarrow I'_{V} = \frac{1}{4} I_{V}$$
$$\mathbf{E}'_{H} = \mathbf{e}_{H} \left( \mathbf{e}_{H} \cdot \mathbf{E}_{D+} \right) = \mathbf{e}_{H} \left( \mathbf{e}_{H} \cdot \mathbf{e}_{D+} \right) \left( \mathbf{e}_{D+} \cdot \mathbf{e}_{V} E_{V} \right) = \frac{1}{2} \mathbf{E}_{V} \Longrightarrow I'_{H} = \frac{1}{4} I_{V}$$



### Repeated measurements don't "commute"



Malus's Law:

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$$\mathbf{E}_{D+} = \mathbf{e}_{D+} \left( \mathbf{e}_{D+} \cdot \mathbf{E}'_{V} \right) = \mathbf{e}_{D+} \left( \mathbf{e}_{D+} \cdot \mathbf{e}_{V} \right) \left( \mathbf{e}_{V} \cdot \mathbf{e}_{V} E_{V} \right) = \mathbf{e}_{D+} \frac{E_{V}}{\sqrt{2}} \Longrightarrow I_{D+} = \frac{1}{2} I_{V}$$
$$\mathbf{E}_{D-} = \mathbf{e}_{D-} \left( \mathbf{e}_{D-} \cdot \mathbf{E}'_{V} \right) = \mathbf{e}_{D-} \left( \mathbf{e}_{D-} \cdot \mathbf{e}_{V} \right) \left( \mathbf{e}_{V} \cdot \mathbf{e}_{V} E_{V} \right) = \mathbf{e}_{D-} \frac{E_{V}}{\sqrt{2}} \Longrightarrow I_{D-} = \frac{1}{2} I_{V}$$



## From wave intensity to photon events

Light beam = "stream of photon"  $I = \Phi \hbar \omega$ 

In the "law of large numbers"

Fraction of intensity in given alternative



 $p_{out}^{(i)} = \frac{I_{out}^{(i)}}{r}$ 







### From Malus's Law to Born's Rule



Malus's Law: 
$$I_V = \frac{1}{2} \mathbf{E}_V^* \cdot \mathbf{E}_V = \left| \mathbf{e}_V \cdot \vec{\varepsilon} \right|^2 \left( \frac{1}{2} \left| E_0 \right|^2 \right) = \cos^2 \theta I_0$$

$$p_V = \frac{I_V}{I_0} = \left| \mathbf{e}_V \cdot \vec{\varepsilon} \right|^2 = \left| \cos \theta \right|^2$$

Generalize: Given photon prepared in polarization  $\vec{\varepsilon}_{in}$ , probability of finding  $\vec{\varepsilon}_{out}$ ,  $p_{out} = |\vec{\varepsilon}_{out}^* \cdot \vec{\varepsilon}_{in}|^2$  Born's Rule



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### From Photon Polarization to "Hilbert Space"

#### Complex vector space of dimension *d*. (Hilbert space can be infinite dimensional)

**Dirac Notation:** 

<u>"Kets" = vectors.</u>  $|V\rangle$  <u>"Bras" = dual vectors.</u>  $\langle V|$ 

<u>Inner product = braket</u>  $\langle V | V \rangle = \| V \|^2$ 

Like  $\vec{\varepsilon}$  and  $\vec{\varepsilon}^*$  for photon polarization,  $\|\vec{\varepsilon}\|^2 = \vec{\varepsilon}^* \cdot \vec{\varepsilon}$ 





### **Matrix Representation**

Orthonormal basis 
$$\{|e_i\rangle|i=1,...,d\}$$
  $\langle e_i|e_j\rangle = \delta_{ij} = \begin{cases} 0 & i \neq j \\ 1 & i=j \end{cases}$   
 $|V\rangle = \sum_i V_i |\mathbf{e}_i\rangle$   $V_i = \langle e_i|V\rangle$ 





$$\langle V | \doteq \begin{bmatrix} V_1^* & \cdots & V_d^* \end{bmatrix} = |V\rangle^\dagger$$

Inner product via matrix multiplication

$$\langle U|V \rangle = \begin{bmatrix} U_1^* & \cdots & U_d^* \end{bmatrix} \begin{bmatrix} V_1 \\ \vdots \\ V_d \end{bmatrix} = \sum_i U_i^* V_i$$





### **Linear Operators**

 $\begin{array}{ll} \underline{\mathsf{Maps}} & \hat{A}|U\rangle = |V\rangle & \hat{A}|W\rangle = |Y\rangle \\ \\ \underline{\mathsf{Linear}} & \hat{A}(a|U\rangle + b|W\rangle) = a\hat{A}|U\rangle + b\hat{A}|V\rangle = a|V\rangle + b|Y\rangle \end{array}$ 

$$\underbrace{\text{Matrix Representation}}_{A \doteq \begin{bmatrix} A_{11} & \dots & A_{1d} \\ \vdots & \ddots & \vdots \\ A_{d1} & \cdots & A_{dd} \end{bmatrix}$$

$$A_{ij} = \left\langle e_i \left| \left( \hat{A} \right| e_j \right\rangle \right) \equiv \left\langle e_i \left| \hat{A} \right| e_j \right\rangle$$

$$\langle e_1 | \hat{A} | e_2 \rangle = \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} A_{12} \\ A_{22} \end{bmatrix} = A_{12}$$

I. H. Deutsch, *University of New Mexico* Short Course in Quantum Information

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# **Special Operator Relations**

<u>Characteristic vectors/values</u> (eigenvectors, eigenvalues)  $\hat{A}|V_a\rangle = a|V_a\rangle$ Real eigenvalues,  $\hat{A}^{\dagger} = \hat{A}$ Hermitian operator: Orthogonal eigenvectors <u>Unitary:</u>  $\hat{U}^{\dagger}\hat{U} = \hat{I}$ Preserves the inner product  $\hat{U}|V\rangle = |V'\rangle$   $\hat{U}|W\rangle = |W'\rangle$  $\langle V'|W'\rangle = \langle V|\hat{U}^{\dagger}\hat{U}|W\rangle = \langle V|W\rangle$ 





# The Postulates of QM (I)

• A "pure state" of the system, representing some maximal knowledge, is a normalized ket in Hilbert space  $|\psi\rangle$ .

 $\langle \boldsymbol{\psi} | \boldsymbol{\psi} \rangle = \left\| \boldsymbol{\psi} \right\|^2 = 1$ 

- To every observable we associate a Hermitian operator  $\hat{A}$ .
- The possible values of an observable that one finds in a measurement are its *eigenvalues*, {*a*}.
- The probability of finding a value *a* is given by the square of the amplitude in that eigenvector.

$$P_{a|\psi} = \left| \left\langle e_a \left| \psi \right\rangle \right|^2$$





# Notes

• The kets  $|\psi\rangle$  and  $|\psi'\rangle = e^{i\phi}|\psi\rangle$  are the same state.

$$P_{a|\psi'} = \left| \left\langle e_a \left| \psi' \right\rangle \right|^2 = \left| \left\langle e_a \left| \psi \right\rangle \right|^2 = P_{a|\psi'}$$

• The *relative phase* in a superposition is *very* important.

$$\left|\psi\right\rangle = r_{1}e^{i\phi_{1}}\left|u_{1}\right\rangle + r_{2}e^{i\phi_{2}}\left|u_{2}\right\rangle$$

$$P_{a|\psi} = \left| \left\langle e_a \left| \psi \right\rangle \right|^2 = r_1^2 \left| \left\langle e_a \left| u_1 \right\rangle \right|^2 + r_2^2 \left| \left\langle e_a \left| u_2 \right\rangle \right|^2 \right. \right. \\ \left. + r_1 r_2 \left( e^{i(\phi_1 - \phi_2)} \left\langle e_a \left| u_1 \right\rangle^* \left\langle e_a \left| u_2 \right\rangle + e^{-i(\phi_1 - \phi_2)} \left\langle e_a \left| u_2 \right\rangle^* \right) \right. \right]$$

 $= p_1 p_{a|1} + p_2 p_{a|2} +$ inteference





# Notes

• The set of possible measurement outcomes for an observable "span" the space \_\_\_\_> complete set.

$$|\psi\rangle = \sum_{a} c_{a} |e_{a}\rangle$$
  $c_{a} = \langle e_{a} |\psi\rangle$   $p_{a|\psi} = |\langle e_{a} |\psi\rangle|^{2} = |c_{a}|^{2}$ 

$$\langle \psi | \psi \rangle = \sum_{a} |c_{a}|^{2} = \sum_{a} p_{a|\psi} = 1$$
 Total probability = 1  
 $\langle A \rangle_{\psi} = \sum_{a} a p_{a|\psi} = \sum_{a} a |\langle e_{a} | \psi \rangle|^{2} = \langle \psi | \hat{A} | \psi \rangle$  Expectation value

Not all operators share the same eigenvectors

*Impossible* to measure all observables simultaneously. Measurement of one observable disturbs the other.

# Heisenberg Uncertainty principle





# Postulates of QM (II)

• After a measurement, the state of the system is updated to the eigenvector associated with the measurement outcome.

$$|\psi\rangle \Rightarrow |e_a\rangle$$
 (Bayesian)

• In the absence of any measurement, a *closed system* evolves according to a unitary map (preserving inner products).

$$\left|\psi(t+\tau)\right\rangle = \hat{U}(\tau)\left|\psi(t)\right\rangle$$

• Physics determines the dynamics  $\frac{d}{dt}\hat{U}(t) = \frac{-i}{\hbar}\hat{H}\hat{U}(t)$ 

# The (time dependent) Schrödinger Equation





### **Example: Photon Polarization**

- State: Normalized complex polarization vector  $\vec{\epsilon}$ ,  $\vec{\epsilon}^* \cdot \vec{\epsilon} = 1$
- Orthonormal bases:  $\{\mathbf{e}_{H},\mathbf{e}_{V}\}$   $\{\mathbf{e}_{D_{+}},\mathbf{e}_{D_{-}}\}$   $\{\mathbf{e}_{R},\mathbf{e}_{L}\}$

Let us call {H,V} the "standard basis"

$$\mathbf{e}_{H} = \begin{bmatrix} 1\\ 0 \end{bmatrix} \quad \mathbf{e}_{V} = \begin{bmatrix} 0\\ 1 \end{bmatrix} \qquad \mathbf{e}_{D+} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\ 1 \end{bmatrix} \quad \mathbf{e}_{D-} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\ -1 \end{bmatrix} \qquad \mathbf{e}_{R} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\ i \end{bmatrix} \quad \mathbf{e}_{L} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\ -i \end{bmatrix}$$

• Three (incompatible) observables:



Pauli matrices with eigenvalues ±1





# **Incomplete Information**

Suppose someone prepares a photon by flipping a coin.







# **Incomplete Information**



$$p_{D+} = p_H p_{D+|H} + p_V p_{D+|V}$$
$$p_{D+} = \frac{1}{2} \frac{1}{2} + \frac{1}{2} \frac{1}{2} = \frac{1}{2}$$

Unpolarized!

#### **Density Operator**





# **Quantum Computing vs. Wave Computing**

- Analog vs. Digital.
  - Quantum phenomena involve *discrete events*
  - (e.g. photon detection).
  - Wave and Particle  $\Leftrightarrow$  Analog and Digital.
  - Possibility of error correction.
- Physical space vs. Hilbert space
  - Modes of field require physical resources.
  - Exponential number of modes not scalable.
  - Quantum mechanics allows exponentially large dimension with polynomial physical resources.
  - Entangled states!





# Summary

- Quantum Mechanics predicts the outcomes of experiments.
- States = Vectors in a complex linear space.
- Observables = Matrices (operators) on the space
  - Eigenvalues = possible measurement outcomes.
  - Eigenvectors = the state after the measurement is done.
- Probability of finding a measurement value = square of complex amplitude (Born rule).
- In a closed system, the system dynamics described by unitary matrices.
- Quantum "coherent superpositions" differ from "statistical mixtures". Noise can "decoher" a system.
- Quantum computing not equivalent to wave computing.



